

A Lattice Spanning-Tree Entropy Function

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Abstract The function

$$W(a, b) = \int_0^{2\pi} dx \int_0^{2\pi} dy \ln[1 - a \cos x - b \cos y - (1 - a - b) \cos(x + y)]$$

which expresses the spanning-tree entropy for various two dimensional lattices, for example, is evaluated directly in terms of standard functions. It is applied to derive several limiting values of the Triangular lattice Green function.

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Introduction

The function

$$W(a, b) = \int_0^{2\pi} dx \int_0^{2\pi} dy \ln[1 - a \cos x - b \cos y - c \cos(x + y)]$$

with $a + b + c = 1$ arises frequently in the statistical physics and combinatorics of two dimensional lattice systems. For example:

(i) $a = b = 1/2$

$$S_{sq} = \frac{\ln(2)}{2\pi^2} + \frac{W(a, b)}{4\pi^2}$$

is the spanning-tree entropy for the square lattice[1].

(ii) $a = b = 1/3$

$$S_{tr} = \frac{\ln(6)}{4\pi^2} + \frac{W(a, b)}{4\pi^2}$$

is the spanning-tree entropy for the triangular lattice[2].

(iii)

$$a = \frac{\sinh K_1}{\sinh K_1 + \sinh K_2 + \sinh K_3}$$

$$b = \frac{\sinh K_2}{\sinh K_1 + \sinh K_2 + \sinh K_3}$$

$$F_I = \ln(2) + \frac{1}{8\pi^2} \ln[\sinh K_1 + \sinh K_2 + \sinh K_3] + \frac{W(a, b)}{8\pi^2}$$

is the critical free energy of the Ising model on a triangular lattice[3].

By comparing the free energies of various Potts models which are known to be related and have been worked out in different ways, Chen and Wu[4] have proposed that

$$\begin{aligned} & \frac{1}{4\pi^2} \int_0^{2\pi} d\theta \int_0^{2\pi} d\phi \ln[A + B + C - A \cos \theta - B \cos \phi - C \cos(\theta + \phi)] \\ &= -\ln(2S) + \frac{2}{\pi} [Ti_2(AS) + Ti_2(BS) + Ti_2(CS)] \end{aligned}$$

where $A, B, C \geq 0$ and $S = 1/\sqrt{AB + BC + CA}$. The aim of this note is to provide a direct proof of this formula.

Calculation

By symmetry, one easily finds

$$W(a, b) = 2[W_+(a, b) + W_-(a, b)] \quad (1)$$

where

$$W_{\pm}(a, b) = \int_0^{\pi} dx \int_0^{\pi} dy \ln[1 - a \cos x - b \cos y - c \cos(x \pm y)].$$

Next, we make the standard change of variable

$$u = \tan(x/2) \quad v = \tan(y/2)$$

to obtain

$$\begin{aligned} W_+(a, b) &= 4 \int_0^\infty \int_0^\infty \frac{dudv}{(1+u^2)(1+v^2)} \ln\left[\frac{2}{(1+u^2)(1+v^2)}\right] \\ &+ 4 \int_0^\infty \int_0^\infty \ln[(u+v)^2 - a((u+v)^2 - u^2(1+v^2)) - b((u+v)^2 - v^2(1+u^2))] \frac{dudv}{(1+u^2)(1+v^2)}. \end{aligned}$$

The first integral is elementary, giving

$$\begin{aligned} W_+(a, b) &= -3\pi^2 \ln(2) \\ &+ 4 \int_0^\infty \int_0^\infty \frac{\ln[c(u+v)^2 + au^2(1+v^2) + bv^2(1+u^2)]}{(1+u^2)(1+v^2)} dudv. \end{aligned}$$

Similarly,

$$\begin{aligned} W_-(a, b) &= -3\pi^2 \ln(2) \\ &+ 4 \int_0^\infty \int_0^\infty \frac{\ln[c(u-v)^2 + au^2(1+v^2) + bv^2(1+u^2)]}{(1+u^2)(1+v^2)} dudv. \end{aligned}$$

By inserting the last two expressions into (1) and noting that the resulting integrand is even in v , we have

$$W(a, b) = -12\pi^2 \ln(2) + 8F(a, b),$$

where

$$F(a, b) = \int_0^\infty \frac{du}{1+u^2} \int_{-\infty}^\infty \frac{dv}{1+v^2} \ln[(1-b)u^2 + (1-a)v^2 + (a+b)u^2v^2 + 2cuv].$$

Next, let $u = vw$ to obtain

$$\begin{aligned} F(a, b) &= 2 \int_0^\infty \frac{udu}{1+u^2} \ln(u) \int_{-\infty}^\infty \frac{dw}{1+u^2w^2} \\ &+ \int_0^\infty \frac{du}{u(1+u^2)} \int_{-\infty}^\infty \frac{dw}{w^2+u^{-2}} \ln[1-b + (1-a)w^2 + (a+b)u^2w^2 + 2cw]. \end{aligned}$$

The first integral vanishes and since [5]

$$\int_0^\infty \frac{dw}{w^2+d^2} \ln[\alpha w^2 + 2\beta w + \gamma] = \frac{\pi}{d} \ln[\alpha d^2 + \gamma + 2d\sqrt{\alpha\gamma - \beta^2}],$$

after the substitution $u = 1/z$,

$$F(a, b) = \pi \int_0^\infty \frac{dz}{1+z^2} \ln[(1-a)z^2 + (1+a) + 2\sqrt{[ac + b(1-b)]z^2 + (1-b)(a+b)}]. \quad (2)$$

Now, let us define

$$A = \frac{y \cot(\theta/2)}{1 + \sqrt{1+y^2}} \quad B = \frac{y \tan(\theta/2)}{1 + \sqrt{1+y^2}}.$$

Then $(A+B)/(1-AB) = y \csc \theta$, so $\tan^{-1}(y \csc \theta) = \tan^{-1} A + \tan^{-1} B$. However,

$$\begin{aligned} \csc \theta \tan^{-1}(y \csc \theta) &= -\frac{d}{d\theta} \int_1^{\csc \theta} \frac{\tan^{-1}(yu)}{\sqrt{u^2-1}} du \\ \csc \theta \tan^{-1} A &= -\frac{d}{d\theta} \int_0^A \frac{\tan^{-1} u}{u} du \\ \csc \theta \tan^{-1} B &= \frac{d}{d\theta} \int_0^B \frac{\tan^{-1} u}{u} du. \end{aligned}$$

Hence, since both sides vanish for $\theta = \pi/2$,

$$\int_1^{\csc \theta} \frac{\tan^{-1}(yu)}{\sqrt{u^2-1}} du = Ti_2(A) - Ti_2(B) \quad (3)$$

where

$$Ti_2(z) = \int_0^z \frac{\tan^{-1} x}{x} dx.$$

With

$$\theta = \csc^{-1} \sqrt{\frac{b^2 - a^2 + 1}{1 - a^2}}, \quad y = \sqrt{a^{-2} - 1}, \quad u = \sqrt{\frac{x^2 - a^2 + 1}{1 - a^2}}$$

in (3) one obtains

$$\begin{aligned} \int_0^b \frac{dx}{\sqrt{x^2 - a^2 + 1}} \tan^{-1} \frac{\sqrt{x^2 - a^2 + 1}}{a} = \\ Ti_2\left(\frac{\sqrt{b^2 + 1 - a^2} + b}{1 + a}\right) - Ti_2\left(\frac{\sqrt{b^2 + 1 - a^2} - b}{1 + a}\right). \end{aligned} \quad (4)$$

We next consider the integral

$$g(a, b) = \int_0^\infty \frac{ds}{s^2 + 1} \ln[\sqrt{b^2(s^2 + 1) + 1} + a]$$

for which it is elementary to determine

$$g(a, 0) = \frac{\pi}{2} \ln(1 + a)$$

$$\frac{\partial}{\partial u}g(a, u) = \frac{\tan^{-1}(\sqrt{u^2 + 1 - a^2}/a)}{\sqrt{u^2 + 1 - a^2}}.$$

Therefore, by integrating over u using (4), we find

$$g(a, b) = \frac{\pi}{2} \ln(1 + a) + Ti_2 \left(\frac{\sqrt{b^2 + 1 - a^2} + b}{1 + a} \right) - Ti_2 \left(\frac{\sqrt{b^2 + 1 - a^2} - b}{1 + a} \right)$$

which is easily transformed into

$$\begin{aligned} \int_0^\infty \frac{\ln[\alpha + \sqrt{\beta^2 x^2 + \gamma^2}]}{x^2 + 1} dx &= \frac{\pi}{2} \ln[\beta + \sqrt{\gamma^2 - \alpha^2}] \\ &+ Ti_2 \left(\frac{\alpha + \sqrt{\gamma^2 - \beta^2}}{\beta + \sqrt{\gamma^2 - \alpha^2}} \right) + Ti_2 \left(\frac{\alpha - \sqrt{\gamma^2 - \beta^2}}{\beta + \sqrt{\gamma^2 - \alpha^2}} \right). \end{aligned} \quad (5)$$

The argument of the logarithm in the integrand of $F(a, b)$ in (2) can be factored: Let $R = \sqrt{[ac + b(1 - b)]x^2 + (1 - b)(a + b)}$; then

$$(1 - a)x^2 + 1 + a + 2R =$$

$$[ac + b(1 - b)]^{-1}[a(1 - a) + 2bc + (1 - a)R][a + R].$$

The integrals resulting from inserting this into (2) are either elementary or can be evaluated by using (5). After some algebraic manipulation, we obtain

$$\begin{aligned} W(a, b) &= 4\pi^2 \ln \frac{d}{2} \\ &+ 8\pi[Ti_2(a/d) + Ti_2(b/d) + Ti_2(c/d)] \end{aligned} \quad (6)$$

where $d = \sqrt{ab + bc + ac}$, which is equivalent to Chen and Wu's formula.

In conclusion, we list a few values of the anisotropic triangular lattice Green function that can be obtained from (6) by differentiation. Here, $\Delta(a, b, c) = a + b + c - a \cos x - b \cos y - c \cos(x + y)$, $d = \sqrt{ab + bc + ca}$.

$$\begin{aligned} \int_0^{2\pi} dx \int_0^{2\pi} dy \frac{\Delta(1, 0, 0)}{\Delta(a, b, c)} &= \\ \frac{4\pi}{d^2} (b + c) \left[\frac{\pi}{2} + \frac{a(b + c) + 2bc}{a(b + c)} \tan^{-1}(a/d) - \tan^{-1}\left(\frac{(b + c)d}{bc + d^2}\right) \right] \\ \int_0^{2\pi} dx \int_0^{2\pi} dy \frac{\Delta(0, 0, 1)}{\Delta(a, b, c)} &= \\ \frac{4\pi}{d^2} (a + b) \left[\frac{\pi}{2} + \frac{(a + b)c + 2ab}{(a + b)c} \tan^{-1}(c/d) - \tan^{-1}\left(\frac{(a + b)d}{ab + d^2}\right) \right] \\ \int_0^{2\pi} dx \int_0^{2\pi} dy \frac{\Delta(1, 1, 1)}{\Delta(a, b, c)} &= \frac{4\pi^2}{d^2} (a + b + c) \end{aligned}$$

$$\begin{aligned}
& + \frac{8\pi}{d^2} \left[\frac{bc - a^2}{a} \tan^{-1}(a/d) + \frac{ac - b^2}{b} \tan^{-1}(b/d) + \frac{ab - c^2}{c} \tan^{-1}(c/d) \right]. \\
& \int_0^{2\pi} dx \int_0^{2\pi} dy \frac{\Delta(-1, 1, 0)}{\Delta(a, b, c)} = \frac{2\pi^2}{d^2} (a - b) \\
& + \frac{4\pi}{d^2} \left[(b-a) \tan^{-1}(c/d) + (a+b + \frac{2c(a+b)}{b}) \tan^{-1}(b/d) - (a+b+2c + \frac{2c(a+b)}{a}) \tan^{-1}(a/d) \right]
\end{aligned}$$

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